Mathematics Chalmers & GU

TMA372/MMG800: Partial Differential Equations, 2014-06-10, 8:30-12:30

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Calculators, formula notes and other subject related material are not allowed.

Each problem gives max 6p. Valid bonus poits will be added to the scores.

Breakings: 3: 15-20p, 4: 21-27p och 5: 28p- For GU studentsG:15-24p, VG: 25p-

For solutions and information about gradings see the couse diary in:

http://www.math.chalmers.se/Math/Grundutb/CTH/tma372/1314/index.html

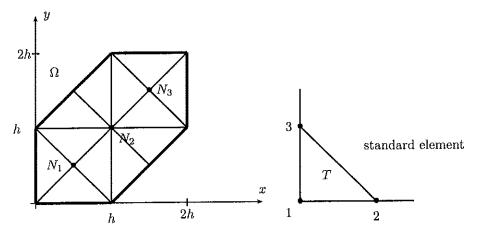
1. Prove the following error estimate for the linear interpolation for a function $f \in C^2(0,1)$,

$$||\pi_1 f - f||_{L_2(a,b)} \le (b-a)^2 ||f''||_{L_2(a,b)}.$$

2. Prove an a priori and an a posteriori error estimate for the cG(1) finite element method for

$$-u''(x) + u'(x) = f$$
, $0 < x < 1$; $u(0) = u(1) = 0$.

3. Let Ω be the hexagonal domain with the uniform triangulation as in the figure below. Compute



the stiffness matrix and the load vector for the cG(1) approximate solution for the problem:

(1)
$$\begin{cases} -\Delta u = 1, & \text{in } \Omega, & \text{(the mesh size is } h) \\ u = 0, & \text{on } \partial\Omega. \end{cases}$$

4. Consider the convection-diffusion problem

$$-\operatorname{div}(\varepsilon\nabla u+\beta u)=f, \text{ in }\Omega\subset\mathbb{R}^2,\quad u=0, \text{ on }\partial\Omega,$$

where Ω is a bounded convex polygonal domain, $\varepsilon > 0$ is constant, $\beta = (\beta_1(x), \beta_2(x))$ and f = f(x). Determine the conditions in the Lax-Milgram theorem that would guarantee existence of a unique solution for this problem. Prove a stability estimate for u i terms of $||f||_{L_2(\Omega)}$, ε and diam (Ω) , and under the conditions that you derived.

5. Consider the boundary value problem

$$(BVP)$$
 $-(a(x)u'(x))' = f(x), \quad 0 < x < 1, \quad u(0) = u(1) = 0, \quad a(x) > 0$

Show that the variational formulation and minimization problem for BVP are equivalent.

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TMA372/MMG800: Partial Differential Equations, 2014-06-10, 8:30-12:30. Solutions.

1. Let $\lambda_0(x) = \frac{\xi_1 - x}{\xi_1 - x_0}$ and $\lambda_1(x) = \frac{x - \xi_0}{\xi_1 - x_0}$ be two linear base functions, where $\xi_0 \neq \xi_1$, ξ_0 , $\xi_1 \in [a, b]$, can be taken as arbitrary interpolation points or just $\xi_0 = a$, $\xi_1 = b$. Then by the integral form of the Taylor formula we may write

$$\begin{cases} f(\xi_0) = f(x) + f'(x)(\xi_0 - x) + \int_x^{\xi_0} (\xi_0 - y) f''(y) \, dy, \\ f(\xi_1) = f(x) + f'(x)(\xi_1 - x) + \int_x^{\xi_1} (\xi_1 - y) f''(y) \, dy, \end{cases}$$

Therefore, the linear function interpolating f in the points $\xi_0, \xi_1 \in [a, b]$, can be written as

$$\Pi_1 f(x) = f(\xi_0) \lambda_0(x) + f(\xi_1) \lambda_1(x)
= f(x) + \lambda_0(x) \int_x^{\xi_0} (\xi_0 - y) f''(y) \, dy + \lambda_1(x) \int_x^{\xi_1} (\xi_1 - y) f''(y) \, dy$$

and by the triangle inequality we get

$$|f(x) - \Pi_{1}f(x)| = \left|\lambda_{0}(x) \int_{x}^{\xi_{0}} (\xi_{0} - y)f''(y) \, dy + \lambda_{1}(x) \int_{x}^{\xi_{1}} (\xi_{1} - y)f''(y) \, dy\right|$$

$$\leq |\lambda_{0}(x)| \left|\int_{x}^{\xi_{0}} (\xi_{0} - y)f''(y) \, dy\right| + |\lambda_{1}(x)| \left|\int_{x}^{\xi_{1}} (\xi_{1} - y)f''(y) \, dy\right|$$

$$\leq |\lambda_{0}(x)| \int_{x}^{\xi_{0}} |\xi_{0} - y||f''(y)| \, dy + |\lambda_{1}(x)| \int_{x}^{\xi_{1}} |\xi_{1} - y||f''(y)| \, dy$$

$$\leq |\lambda_{0}(x)| \int_{x}^{\xi_{0}} (b - a)|f''(y)| \, dy + |\lambda_{1}(x)| \int_{x}^{\xi_{1}} (b - a)|f''(y)| \, dy$$

$$\leq (b - a) \left(|\lambda_{0}(x)| + |\lambda_{1}(x)|\right) \int_{a}^{b} |f''(y)| \, dy$$

$$= (b - a) \left(\lambda_{0}(x) + \lambda_{1}(x)\right) \int_{a}^{b} |f''(y)| \, dy = (b - a) \int_{a}^{b} |f''(y)| \, dy.$$

Through repeated use of the Cauchy's inequality it follows that

$$\int_{a}^{b} |f(x) - \Pi_{1}f(x)|^{2} dx \leq \int_{a}^{b} (b - a)^{2} \left(\int_{a}^{b} |f''(y)| \, dy \right)^{2} dx$$

$$= (b - a)^{3} \left(\int_{a}^{b} |f''(y)| \, dy \right)^{2} = (b - a)^{3} \left(\int_{a}^{b} 1 \cdot |f''(y)| \, dy \right)^{2}$$

$$\leq (b - a)^{3} \int_{a}^{b} 1^{2} \, dy \cdot \int_{a}^{b} |f''(y)|^{2} \, dy$$

$$= (b - a)^{4} \int_{a}^{b} |f''(y)|^{2} \, dy.$$

Consequently

$$\Big(\int_a^b |f(x) - \Pi_1 f(x)|^2 dx\Big)^{1/2} \le (b - a)^2 \Big(\int_a^b |f''(y)|^2 dy\Big)^{1/2},$$

and we have the desired result viz,

$$||f - \Pi_1 f||_{L_2(a,b)} \le (b-a)^2 ||f''||_{L_2(a,b)}$$

2. We multiply the differential equation by a test function $v \in H_0^1(I)$, I = (0,1) and integrate over I. Using integration by parts integration and the boundary conditions we get the following variational problem: Find $u \in H_0^1(I)$ such that

(2)
$$\int_{I} (u'v' + u'v) = \int_{I} fv, \quad \forall v \in H_0^1(I).$$

A Finite Element Method with cG(1) reads as follows: Find $U \in V_h^0$ such that

(3)
$$\int_{I} (U'v' + U'v) = \int_{I} fv, \quad \forall v \in V_h^0 \subset H_0^1(I),$$

where $V_h^0 = \{v : v \text{ is piecewise linear and continuous in a partition } \mathcal{T}_h \text{ of } I, \ v(0) = v(1) = 0\}.$ Now let e = u - U, then (2)–(3) gives that

(4)
$$\int_{I} (e'v' + e'v) = 0, \quad \forall v \in V_h^0.$$

We note that using e(0) = e(1) = 0, we get

(5)
$$\int_{I} e'e = \int_{I} \frac{1}{2} \frac{d}{dx} \left(e^{2} \right) = \frac{1}{2} (e^{2})|_{0}^{1} = 0.$$

Further, using Poincare inequality we have

$$||e||^2 \le ||e'||^2$$
.

A priori error estimate: We use Poincare inequality, Galerkin orthogonality (4), (5) and standard interpolation estimates to get

$$\begin{aligned} \|e\|_{H^{1}}^{2} &= \int_{I} (e'e' + ee) \leq 2 \int_{I} e'e' = 2 \int_{I} (e'e' + e'e) = 2 \int_{I} \left(e'(u - U)' + e'(u - U) \right) \\ &= 2 \int_{I} \left(e'(u - \pi_{h}u)' + e'(u - \pi_{h}u) \right) + 2 \int_{I} \left(e'(\pi_{h}u - U)' + e'(\pi_{h}u - U) \right) \\ &= \{ v = U - \pi_{h}u \text{ in } (4) \} = 2 \int_{I} \left(e'(u - \pi_{h}u)' + e'(u - \pi_{h}u) \right) \\ &\leq 2 \|(u - \pi_{h}u)'\| \|e'\| + 2 \|u - \pi_{h}u\| \|e'\| \\ &\leq 2 C_{i} \{ \|hu''\| + \|h^{2}u''\| \} \|e\|_{H^{1}}, \end{aligned}$$

this gives that

$$||e||_{H^1} \le C_i \{||hu''|| + ||h^2u''||\}.$$

which is the a priori error estimate.

A posteriori error estimate:

$$||e||_{H^{1}}^{2} = \int_{I} (e'e' + ee) \leq 2 \int_{I} e'e' = 2 \int_{I} (e'e' + e'e)$$

$$= 2 \int_{I} ((u - U)'e' + (u - U)'e) = \{v = e \text{ in } (4)\}$$

$$= 2 \int_{I} fe - \int_{I} (U'e' + U'e) = \{v = \pi_{h}e \text{ in } (5)\}$$

$$= \int_{I} f(e - \pi_{h}e) - \int_{I} \left(U'(e - \pi_{h}e)' + U'(e - \pi_{h}e)\right)$$

$$= \{P.I. \text{ on each subinterval}\} = \int_{I} \mathcal{R}(U)(e - \pi_{h}e),$$

where $\mathcal{R}(U) := f + U'' - U' = f - U'$, (for approximation with piecewise linears, $U'' \equiv 0$, on each subinterval). Thus Cauchy Schwars and standard interpolation estimates implies that

$$\|e\|_{H^1}^2 \leq \|h\mathcal{R}(U)\|\|h^{-1}(e-\pi_h e)\| \leq C_i \|h\mathcal{R}(U)\|\|e'\| \leq C_i \|h\mathcal{R}(U)\|\|e\|_{H^1},$$

where C_i is an interpolation constant, and hence we have with $\|\cdot\| = \|\cdot\|_{L_2(I)}$ that

$$||e||_{H^1} \leq C_i ||h\mathcal{R}(U)||.$$

3. Let V be the linear function space defined by

$$V := \{v : v \in H^1(\Omega), v = 0, \text{ on } \partial\Omega\}.$$

Multiplying the differential equation by $v \in V$ and integrating over Ω we get that

$$-(\Delta u, v) = (1, v), \quad \forall v \in V.$$

Now using Green's formula and the fact that v = 0 on $\partial \Omega \setminus \Gamma_1$, we have that

$$-(\Delta u, \nabla v) == (\nabla u, \nabla v) - \int_{\partial \Omega} (n \cdot \nabla u) v \, ds = (\nabla u, \nabla v), \qquad \forall v \in V,$$

Hence, the variational formulation is:

$$(\nabla u, \nabla v) = (1, v), \quad \forall v \in V$$

Let V_h be the usual finite element space consisting of continuous piecewise linear functions satisfying the boundary condition v=0 on $\partial\Omega$: Then, the cG(1) method is: Find $U\in V_h$ such

$$(\nabla U, \nabla v) = (1, v) \quad \forall v \in V_h$$

 $(\nabla U, \nabla v) = (1, v) \qquad \forall v \in V_h$ Making the "Ansatz" $U(x) = \sum_{j=1}^3 \xi_j \varphi_j(x)$, where $\varphi_j, j = 1, 2, 3$ are the standard basis functions corresponding to the interior nodes N_1, N_2 and N_3 , we obtain the system of equations

$$\sum_{i=1}^{3} \xi_{j} \int_{\Omega} \nabla \varphi_{i} \cdot \nabla \varphi_{j} \, dx = \int_{\Omega} f \varphi_{i} \, dx, \quad i = 1, 2, 3.$$

In matrix form this can be written as $S\xi = F$, where $S_{ij} = (\nabla \varphi_i, \nabla \varphi_j)$ is the stiffness matrix, and $F_i = (f, \varphi_i)$ is the load vector.

We first compute the stiffness matrix for the reference triangle T. The local basis functions are

$$\phi_{1}(x_{1}, x_{2}) = 1 - \frac{x_{1}}{h} - \frac{x_{2}}{h}, \qquad \nabla \phi_{1}(x_{1}, x_{2}) = -\frac{1}{h} \begin{bmatrix} 1 \\ 1 \end{bmatrix},$$

$$\phi_{2}(x_{1}, x_{2}) = \frac{x_{1}}{h}, \qquad \nabla \phi_{2}(x_{1}, x_{2}) = \frac{1}{h} \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$

$$\phi_{3}(x_{1}, x_{2}) = \frac{x_{2}}{h}, \qquad \nabla \phi_{3}(x_{1}, x_{2}) = \frac{1}{h} \begin{bmatrix} 0 \\ 1 \end{bmatrix}.$$

Hence, with $|T| = \int_T dz = h^2/2$, we can easily compute

$$\begin{split} s_{11} &= (\nabla \phi_1, \nabla \phi_1) = \int_T |\nabla \phi_1|^2 \, dx = \frac{2}{h^2} |T| = 1, \\ s_{12} &= s_{21} = (\nabla \phi_1, \nabla \phi_2) = \int_T \frac{-1}{h^2} |T| = -1/2, \\ s_{23} &= s_{32} = (\nabla \phi_2, \nabla \phi_3) = 0, \\ s_{22} &= s_{33} = \dots = \frac{1}{h^2} |T| = 1/2. \end{split}$$

Thus by symmetry we get that the local stiffness matrix for the standard element is:

$$s = \frac{1}{2} \left[\begin{array}{ccc} 2 & -1 & -1 \\ -1 & 1 & 0 \\ -1 & 0 & 1 \end{array} \right].$$

We can now assemble the global stiffness matrix S from the local stiffness matrix s:

$$S_{11} = 4s_{11} = 4,$$
 $S_{12} = S_{21} = 2s_{12} = -1$ $S_{13} = s_{31} = 0$ $S_{22} = 8s_{22} = 4$ $S_{23} = 2s_{12} = -1,$ $S_{33} = 4s_{11} = 4.$

The remaining matrix elements are obtained by symmetry $S_{ij} = S_{ji}$. Hence,

$$S = \left[\begin{array}{rrr} 4 & -1 & 0 \\ -1 & 4 & -1 \\ 0 & -1 & 4 \end{array} \right].$$

As for the load vector we have that

$$\int_{\Omega} \varphi_1 = \int_{\Omega} \varphi_3 = 4\frac{1}{3}\frac{h^2}{2}.1 = \frac{2}{3}h^2, \quad \int_{\Omega} \varphi_2 = 8\frac{1}{3}\frac{h^2}{2}.1 = \frac{4}{3}h^2.$$

This the load vector is given by $b = \frac{h^2}{3}(2,4,2)^t$. Observe that, here S has become independent of h.

4. Consider

(7)
$$-\operatorname{div}(\varepsilon \nabla u + \beta u) = f, \text{ in } \Omega, \qquad u = 0 \text{ on } \Gamma = \partial \Omega.$$

a) Multiply the equation (7) by $v \in H_0^1(\Omega)$ and integrate over Ω to obtain the Green's formula

$$-\int_{\Omega} \operatorname{div}(\varepsilon \nabla u + \beta u) v \, dx = \int_{\Omega} (\varepsilon \nabla u + \beta u) \cdot \nabla v \, dx = \int_{\Omega} f v \, dx.$$

The variational formulation for our problem is now: Find $u \in H_0^1(\Omega)$ such that

(8)
$$a(u,v) = L(v), \qquad \forall v \in H_0^1(\Omega)$$

where

$$a(u, v) = \int_{\Omega} (\varepsilon \nabla u + \beta u) \cdot \nabla v \, dx,$$

and

$$L(v) = \int_{\Omega} fv \, dx.$$

According to the Lax-Milgram theorem, for a unique solution for (7) we need to verify that the following relations are valid:

i)

$$|a(v,w)| \le \gamma ||u||_{H^1(\Omega)} ||w||_{H^1(\Omega)}, \quad \forall v, w \in H^1_0(\Omega),$$

ii)

$$a(v,v) \ge \alpha ||v||_{H^1(\Omega)}^2, \quad \forall v \in H^1_0(\Omega),$$

iii)

$$|L(v)| \le \Lambda ||v||_{H^1(\Omega)}, \qquad \forall v \in H^1_0(\Omega),$$

for some γ , α , $\Lambda > 0$.

Now since

$$|L(v)| = |\int_{\Omega} fv \, dx| \le ||f||_{L_2(\Omega)} ||v||_{L_2(\Omega)} \le ||f||_{L_2(\Omega)} ||v||_{H^1(\Omega)},$$

thus iii) follows with $\Lambda = ||f||_{L_2(\Omega)}$.

Further we have that

$$\begin{split} |a(v,w)| &\leq \int_{\Omega} |\varepsilon \nabla v + \beta v| |\nabla w| \ dx \leq \int_{\Omega} (\varepsilon |\nabla v| + |\beta| |v|) |\nabla w| \ dx \\ &\leq \Big(\int_{\Omega} (\varepsilon |\nabla v| + |\beta| |v|)^2 \ dx \Big)^{1/2} \Big(\int_{\Omega} |\nabla w|^2 \ dx \Big)^{1/2} \\ &\leq \sqrt{2} \max(\varepsilon, ||\beta||_{\infty}) \Big(\int_{\Omega} (|\nabla v|^2 + v^2) \ dx \Big)^{1/2} ||w||_{H^1(\Omega)} \\ &= \gamma ||v||_{H^1(\Omega)} ||w||_{H^1(\Omega)}, \end{split}$$

which, with $\gamma = \sqrt{2} \max(\epsilon, ||\beta||_{\infty})$, gives i).

Finally, if $div\beta \leq 0$, then

$$a(v,v) = \int_{\Omega} \left(\varepsilon |\nabla v|^2 + (\beta \cdot \nabla v)v \right) dx = \int_{\Omega} \left(\varepsilon |\nabla v|^2 + (\beta_1 \frac{\partial v}{\partial x_1} + \beta_2 \frac{\partial v}{\partial x_2})v \right) dx$$

$$= \int_{\Omega} \left(\varepsilon |\nabla v|^2 + \frac{1}{2} (\beta_1 \frac{\partial}{\partial x_1} (v)^2 + \beta_2 \frac{\partial}{\partial x_2} (v)^2) \right) dx = \text{Green's formula}$$

$$= \int_{\Omega} \left(\varepsilon |\nabla v|^2 - \frac{1}{2} (\operatorname{div}\beta)v^2 \right) dx \ge \int_{\Omega} \varepsilon |\nabla v|^2 dx.$$

Now by the Poincare's inequality

$$\int_{\Omega} |\nabla v|^2 \, dx \geq C \int_{\Omega} (|\nabla v|^2 + v^2) \, dx = C ||v||_{H^1(\Omega)}^2,$$

for some constant $C = C(\Omega)$, we have

$$a(v,v) \geq \alpha ||v||_{H^1(\Omega)}^2, \qquad \text{with } \alpha = C\varepsilon,$$

thus ii) is valid under the condition that $\operatorname{div}\beta \leq 0$.

From ii), (8) (with v = u) and iii) we get that

$$\alpha||u||_{H^1(\Omega)}^2 \leq a(u,u) = L(u) \leq \Lambda ||u||_{H^1(\Omega)},$$

which gives the stability estimate

$$||u||_{H^1(\Omega)} \leq \frac{\Lambda}{\alpha},$$

with $\Lambda = ||f||_{L_2(\Omega)}$ and $\alpha = C\varepsilon$ defined above.

To summarize: The conditions for the Lax-Milgram theorem are:

$$f \in L_2(\Omega), \ \beta \in L_{\infty}(\Omega) \ \text{and} \ \nabla \cdot \beta \leq 0$$
 a.e.

5. See the Book and/or Lecture Notes.

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