#### Quantum Mechanics FKA081/FIM400

#### Final Exam 19 August 2014

Next review time for the exam: September 26 any time (9-18) in my room.

NB: If you want to come to the review you must collect your exam before at the "Kansli" in Origo 5th floor. (This info is also available on the course homepage.)

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#### Allowed material during the exam:

- The course textbook J.J. Sakurai and Jim Napolitano, Modern Quantum Mechanics Second Edition (2010).
   NB: The old red cover version: J.J. Sakurai, Modern Quantum Mechanics Revised Edition (1994) is also allowed.
- A Chalmers approved calculator.

Write the final answers clearly marked by Ans: ... and underline them.

You may use without proof any formula in the book.

The grades are assigned according to the table in the course homepage.

# Problem 1

Consider a spinless particle in one dimension in an arbitrary state  $|\psi\rangle$  and denote by x and p the position and momentum operator respectively.

Q1 (1 points) Show that  $\langle \psi | xp | \psi \rangle$  can never be equal to  $\langle \psi | px | \psi \rangle$ .

Q2 (2 points) Show that neither  $\langle \psi | xp | \psi \rangle$  nor  $\langle \psi | px | \psi \rangle$  can ever be real. (In particular, they can never be zero.)

Q3 (2 points) Find the linear combination of  $\langle \psi | x p | \psi \rangle$  and  $\langle \psi | p x | \psi \rangle$  that is always real and give an example where this combination vanishes.

# Problem 2

Consider the following hamiltonian for a two state system (units of eV understood, no need to write them.)

$$H_1 = \begin{pmatrix} 1 & a \\ a & 2 \end{pmatrix} \tag{1}$$

Q1 (2 points) For generic small a make a graph comparing the exact eigenvalues with second order in perturbation theory

Consider now instead

$$H_2 = \begin{pmatrix} 1 & a \\ a & 1 \end{pmatrix} \tag{2}$$

Q2 (3 points) Discuss the most important differences that arise in this case and the reason for them.

## Problem 3

Consider two states  $|\psi_0\rangle$  and  $|\psi_E\rangle$  of energy 0 and E > 0 respectively. At times t < 0 the system is in the ground state  $|\psi_0\rangle$ . Between time t = 0 and time t = T, we turn on a perturbation H' mixing the two states, i.e.

$$\langle \psi_0 | H' | \psi_E \rangle = \langle \psi_E | H' | \psi_0 \rangle = \epsilon \ll E, \quad \langle \psi_0 | H' | \psi_0 \rangle = \langle \psi_E | H' | \psi_E \rangle = 0.$$
(3)

**Q1** (2 points) Use first order perturbation theory to find the probability that the system is in the excited state after t > T.

Now consider the same system with the following experimental values:

$$\Omega = \frac{E}{\hbar} = 1.00 \times 10^{12} \text{ Hz}, \quad \omega = \frac{\epsilon}{\hbar} = 2.00 \times 10^{11} \text{ Hz}, \quad T = 3.00 \times 10^{-12} \text{ s}$$

Q2 (1 points) Give the numerical answer to probability found in Q1.

Q3 (4 points) Solve the problem numerically with the above values without using perturbation theory and check how well the result agrees with perturbation theory.

(*Hints:* You never need to use the value of  $\hbar$  if you always work with frequencies. Consider the exact total hamiltonian  $(1/\hbar)H$  in the interval between 0 and T. It's OK if you get negative eigenfrequencies, they just arise because we chose the energy of the unperturbed ground state to be zero.)

### Problem 4

Consider a nucleus of spin s = 5/2 subjected to the following Hamiltonian: ( $\gamma$  and  $\beta$  are constants)

$$H = \left(\beta S_x^2 + \gamma (S_y^2 + S_z^2)\right) \tag{4}$$

Q1 (2 points) Find the spectrum of H.

Q2 (2 points) Find the degeneracy of the eigenvalues in the spectrum of H.

## Problem 5

Consider the two observables:

$$A = \begin{pmatrix} 1 & 1 \\ 1 & 2 \end{pmatrix}, \qquad B = \begin{pmatrix} 2 & x \\ x & 3 \end{pmatrix}$$
(5)

**Q1** (1 point) For what value of x can they be simultaneously diagonalized? **Q2** (1 point) Find the eigenvalues of A and B for x found in **Q1**.

Q3 (2 points) How are the eigenvalues found in Q2 paired? In other words, which pairs correspond to the same eigenvector?

# Problem 6

In the ground state of a hydrogen atom, the electron is subjected at t > 0 to the perturbation

$$H' = gr\cos\theta e^{-\gamma t} \tag{6}$$

where r and  $\theta$  are the usual polar coordinates of the electron, g is a small parameter and  $\gamma > 0$ . You may ignore the spin degree of freedom in this problem and use lowest order perturbation theory.

**Q1** (2 points) Which of the states  $|n, l, m\rangle$  can be excited by this perturbation?

Q2 (3 points) What is the probability for the electron to be found in the first excited state (n = 2) state at  $t = +\infty$ ?

You may use the fact that

$$\int_{0}^{\infty} r^{3} \mathcal{R}_{21}(r) \mathcal{R}_{10}(r) \, \mathrm{d}r = \frac{2^{7} \sqrt{2}}{3^{4} \sqrt{3}} a_{0},$$
$$\langle Y_{10} | \cos \theta | Y_{00} \rangle = \frac{1}{\sqrt{3}}$$

where  $a_0$  is the Bohr radius and  $\mathcal{R}_{nl}(r)$  are the radial wave functions.

PROBLEM 1

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Q1:  $\langle \psi | x p | \psi \rangle - \langle \psi | p x | \psi \rangle =$ =<4[[x,p]14> = ifi<414>=ifi=0 Q2:  $\langle \psi | \times p | \psi \rangle^* = \langle \psi | (x p)^{\dagger} | \psi \rangle =$  $= \langle \psi | P^{\dagger} \times^{\dagger} | \psi \rangle = \langle \psi | P \times | \psi \rangle$ =>  $Re(\langle \psi | x p | \psi \rangle) = Re(\langle \psi | p \times | \psi \rangle)$  $I_m(\langle \psi | x p | \psi \rangle) = - I_m(\langle \psi | p \times | \psi \rangle)$ By Q1: 2 Jm (<+1×p1+>)=iti=0. Q3: ByQ2: <41xp14>+<41px14>eR on a real wave function;  $\langle \psi | x \rho | \psi \rangle = -i \int dx \psi (x) x \psi'(x)$ purely maginenz. Same for (41PXH) Thus the sum (which is real) must be zero.

PROBLEM 2  
Q1. Exact Sol: 
$$\begin{vmatrix} 1-E & q \\ a & 2-E \end{vmatrix} = 0$$
  
 $\Rightarrow E^{2} - 3E + 2a^{2} = 0 \Rightarrow E_{1/2} = \frac{3\pm\sqrt{1+4a^{2}}}{2}$   
Pert. Heavy:  $H_{0} = \begin{pmatrix} 0 & 2 \\ 0 & 2 \end{pmatrix} + H'_{2} \begin{pmatrix} 0 & q \\ q & 0 \end{pmatrix}$   
 $HY = \begin{pmatrix} 0 \\ 0 \end{pmatrix} + 12Y = \begin{pmatrix} 0 \\ 1 \end{pmatrix} + 11Y = \begin{pmatrix} 0 & q \\ q & 0 \end{pmatrix}$   
 $I_{0} 2062: \langle 1|H'_{11}Y \rangle = \langle 2|H'_{12}Y \rangle \equiv 0$ .  
I order:  $\Delta E_{(1)}^{\pi} = \frac{a}{1-2} = -a^{2}$   
 $E = \frac{2-1}{2-1} = 2+a^{2}$   
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Q2: Exact sol: | 1-E9 |= 0 =>  $E^{2} - 2E + 1 - a^{2} = 0 => E_{1,2} = 1 \pm a$ The exact sol. is the same as the one that is found in degenerate first order pert. theory:  $H_{0} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad H' = \begin{pmatrix} 0 & q \\ q & 0 \end{pmatrix}$ (eigenvalues of H': ±a). 1+9 1-9 9 Rifferences between QI and QZ: .) Non-degenerate us. degenerate. (AE a a DE a a) 000) Approx U.S exact  $[H_{\varrho}, H'] = O$ . [Ho, H'] = O

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$$\frac{P_{ROBLSM}}{P_{ROBLSM}} = \frac{i}{t} \int_{0}^{t} \langle \Psi_{E} | V(t) | \Psi_{O} \rangle e^{i \frac{t}{t}} dt$$

$$\frac{i Et}{t} dt$$

$$\int_{excite} \int_{excite} \int_{excite$$

Q2  
Numerically  
Pexcile = 
$$\frac{4 \times (0,2)}{1^2} \sin^2(\frac{1\cdot 3}{2}) = 0,159$$

Q3 For 
$$t \in [0, T]$$
 the  
hamiltonian is:  
 $\frac{1}{2}H = \begin{pmatrix} 0 & e \\ e & E \end{pmatrix} = \begin{pmatrix} 0 & 0, 2 \\ 0, 2 & 1 \end{pmatrix} * 10^{12} Hz.$   
eigenvectors:  $\phi = \begin{pmatrix} 0, 189 \\ 0, 982 \end{pmatrix}$  with  $W_1 = 1, 04 \cdot 10^{12}$   
and  $\phi_2 = \begin{pmatrix} -0, 982 \\ 0, 189 \end{pmatrix}$  with  $W_2 = -3,85 \cdot 10^{12}$   
The normalized solution that  
matches  $\Psi_0 = \begin{pmatrix} 0 \end{pmatrix}$  when  $t = 0$  is:  
 $\Psi(t) = 0, 189 \phi_1 e^{-1} = -3,85 \cdot 10^{12}$   
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 $\Psi(t) = 0, 189 \phi_1 e^{-1} = -3,85 \cdot 10^{12}$   
 $\Psi(t) = 0, 189 \phi_1 e^{-1} = 0,138$   
The perturbative result is  $\approx 15\%$  off.

PROBLEM 4. Q1/2: Write  $H = \gamma S^2 + (\beta - \gamma)S_x^2$ where  $S_{-}^2 S_x^2 + S_y^2 + S_z^2$  $[S^2, S^2_x] = 0$  hence they can be simultaneously die genelized. The possible eigenvalues of Sx are the same as those of Sz (Just relable the exes). Hence the spectrum is:  $E = \gamma_{3(3+1)} + (\beta - \gamma) m^{2}$ where S = 5/2 and m = - 5/3/2/2/5  $E = \begin{cases} \frac{5}{2}\chi + \frac{25}{4}\beta & m = \pm \frac{5}{2} & doubly deg \\ \frac{13\chi}{2} + \frac{9\beta}{4} & m = \pm \frac{3}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{\beta}{4} & m = \pm \frac{12}{2} & 1111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 1111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 1111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 1111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 1111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 1111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 11111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 11111 \\ \frac{17\chi}{2} + \frac{12\chi}{4} & m = \pm \frac{12\chi}{2} & 11111 \\ \frac{111}{2} + \frac{12\chi}{2} & \frac{12\chi}{4} & \frac$ m=±5 doubly deg.

$$P_{ROBLSM5}$$

$$QI: [A, B] = \begin{pmatrix} 0 & 1-x \\ -1+x & 0 \end{pmatrix} = 0 \iff x=1.$$

Q2:  $det(A - q1) = a^2 - 3a + 1 = 0$ => Q1,2= 3±15  $det(B-b4)=b^2-5b+5=0$ => b12 = 5±15 Q3: The easiest way is to find an eigenvector of A:  $\left(1 - \frac{3+1}{2} + \frac{1}{2} + \frac{1}{2}\right) \left[\frac{\alpha}{\beta}\right] = 0$ AU=  $\alpha_1 U = 2$  (1)  $2 - \frac{3+1}{2} + \frac{1}{2} + \frac{1}{2}$ => X=1, P= 1+V5 up to unimportant normaliz.  $Bv = \begin{pmatrix} 2 \\ 1 \\ 3 \end{pmatrix} \begin{pmatrix} 1 \\ 1+V_5 \\ -2 \end{pmatrix} = \begin{pmatrix} 5+V_5 \\ -2 \\ 5+3V_5 \end{pmatrix} = \frac{5+V_5}{2} \cdot U$ => Pairs:  $\left(\frac{3+\sqrt{5}}{2}, \frac{5+\sqrt{5}}{2}\right)$  and  $\left(\frac{3-\sqrt{5}}{2}, \frac{5-\sqrt{5}}{2}\right)$ 

PROBLEM 6 < 4 1 H' 1 4 1,0,0 > ~ < Yem (cos 0 1 %) QI: => only m=0, l=1 contribute (You can see this by noticing that To=const. and cost x Yio and using I of Tem)  $\frac{\sqrt{2}}{P_{1\to 2}} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{210}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{210} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{100} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{100} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{100} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{100} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{100} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{100} | H^{i} | \Psi_{100} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{10} | \Psi_{10} | \Psi_{10} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{10} | \Psi_{10} | \Psi_{10} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{10} | \Psi_{10} | \Psi_{10} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{10} | \Psi_{10} | \Psi_{10} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{10} | \Psi_{10} | \Psi_{10} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_{21}t} \langle \Psi_{10} | \Psi_{10} | \Psi_{10} \rangle \right|^{2} = \frac{1}{h^2} \left| \int_{0}^{+\infty} dt \, e^{-i W_$  $=\frac{1}{t^2} \int dt e ge \int dr r R(r) \cdot r \cdot R(r) \times \frac{1}{2} \int dt e ge \int dr r R(r) \cdot r \cdot R(r) \times \frac{1}{2} \int dt e ge \int dr r \cdot R(r) \cdot r \cdot R(r) \cdot r \cdot R(r) \times \frac{1}{2} \int dt e ge \int dr r \cdot R(r) \cdot \cdot R(r$  $\times \langle \gamma_{10} | 6 S \theta | \gamma_{00} \rangle |^{2} =$  $=\frac{9}{7}\left|\frac{1}{1-1}\right|^{2}\cdot\left(\frac{2^{1}\sqrt{2}}{3^{4}\sqrt{3}}a_{0}\right)^{2}\cdot\left(\frac{1}{\sqrt{3}}\right)=$  $= \frac{2^{15}}{3^{10}} \cdot \frac{9^{2}}{4^{2}} \cdot \frac{1}{8^{2}} \cdot \frac{9^{2}}{8^{2}} \cdot \frac{1}{8^{2}} \cdot \frac{9^{2}}{8^{2}} \sqrt{(\omega_{21} = \frac{3e^{2}}{8a_{0}})}$