

Quantum Mechanics FKA081/FIM400

Final Exam October 23 2013

Next review time for the exam: Thursday November 21 between 12-13 in my room.

NB: If you want to come to the review you must collect your exam before at the “Kansli” in Origo 5th floor. (This info is also available on the course homepage.)

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Allowed material during the exam:

- The course textbook J.J. Sakurai and Jim Napolitano, Modern Quantum Mechanics Second Edition (2010).
NB: The old red cover version: J.J. Sakurai, Modern Quantum Mechanics Revised Edition (1994) is *also* allowed.
- A Chalmers approved calculator.

Write the final answers clearly marked by **Ans: ...** and underline them.

You may use without proof any formula in the book.

There is a total of 30 points in this test. The grades are assigned according to the table in the course homepage.

Problem 1

Consider the following observable ($\epsilon \ll 1$)

$$A = \begin{pmatrix} 1 & \epsilon & \epsilon \\ \epsilon & 2 & \epsilon \\ \epsilon & \epsilon & 3 \end{pmatrix} \quad (1)$$

Q1 (*1 points*) Compute the eigenvalues to first order in ϵ .

Q2 (*2 points*) Compute the eigenvalues to second order in ϵ .

Problem 2

Consider two identical spin-1/2 particles of mass M , moving in one dimension and interacting with a two-body potential of the form

$$V = \frac{k}{2}(x_1 - x_2)^2 + \alpha(x_1 - x_2)^4 \mathbf{S}_1 \cdot \mathbf{S}_2 \quad (2)$$

Q1 (*3 points*) Compute the energies and the degeneracies of the two lowest states in the center of mass frame for the case $\alpha = 0$.

Q2 (*4 points*) Compute the first order correction in α of the two lowest states.

Problem 3

Consider a system containing three spin-1 particles.

Q1 (*4 points*) Assuming that all three particles are *distinguishable*, write all possible values for the total spin of the system, including possible degeneracies. (You *do not* have to write down the states explicitly).

Q2 (*3 points*) Assuming that all three particles are *indistinguishable*, and that their orbital wave-function is totally antisymmetric, write the possible values for the total spin of the system and the corresponding state explicitly.

Problem 4

Consider a spinless particle in one dimension with a potential $\lambda\delta(x)$ ($\lambda > 0$).

Q1 (*3 points*) Compute the transmission and reflection coefficients.

Problem 5

Consider a spinless particle in a state of orbital angular momentum $l = 1$. (That is, $L_z|\pm\rangle = \pm|\pm\rangle$ and $L_z|0\rangle = 0$.)

The Hamiltonian for the so called quadrupole moment can be described by

$$H = \hbar\omega(L_xL_z + L_zL_x) \quad (3)$$

Q1 (3 points) Find the eigenvalues of the Hamiltonian.

Q2 (4 points) Assuming that at time $t = 0$ the particle is in the state

$$|\psi, t = 0\rangle = \frac{1}{\sqrt{2}}(|+\rangle - |-\rangle) \quad (4)$$

find the state at time $t > 0$.

Problem 6

Consider the one dimensional Hamiltonian ($k > 0$)

$$H = \frac{p^2}{2m} + k|x| \quad (5)$$

Q1 (2 points) Compute the ground state energy using two possible trial wave functions: ($a > 0$)

$$\begin{array}{ll} \text{case 1} & \exp(-a|x|) \\ \text{case 2} & \exp(-ax^2/2) \end{array}$$

Q2 (1 point) Which of the two is better and why?

PROBLEM 1:

Unperturbed states: $|1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, |2\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, |3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$

Q1: Let $A' = \begin{pmatrix} 0 & \epsilon & \epsilon \\ \epsilon & 0 & \epsilon \\ \epsilon & \epsilon & 0 \end{pmatrix} = \epsilon \begin{pmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{pmatrix}$

First order: $\langle 1|A'|1\rangle = \langle 2|A'|2\rangle = \langle 3|A'|3\rangle = 0$ for all.

Q2: First notice that

$$A'_{ij} = \langle i|A'|j\rangle = \epsilon \text{ for } i \neq j.$$

$$\text{Then: } E_1^{(2)} = \sum_{k \neq 1} \frac{|A'_{1k}|^2}{1-k} = \frac{\epsilon^2}{-1} + \frac{\epsilon^2}{-2} = -\frac{3\epsilon^2}{2}$$

$$E_2^{(2)} = \sum_{k \neq 2} \frac{|A'_{2k}|^2}{2-k} = \frac{\epsilon^2}{1} + \frac{\epsilon^2}{-1} = 0$$

$$E_3^{(2)} = \sum_{k \neq 3} \frac{|A'_{3k}|^2}{3-k} = \frac{\epsilon^2}{2} + \frac{\epsilon^2}{1} = \frac{3\epsilon^2}{2}$$

PROBLEM 2

Set $x = x_1 - x_2$, $m = \frac{M}{2}$ (reduced mass).

Q1. For $\alpha=0$ $H_0 = \frac{p^2}{2m} + \frac{k}{2} x^2$

eigenvalues $E_n = \hbar \omega \left(n + \frac{1}{2}\right)$ $\omega = \sqrt{\frac{k}{m}} = \sqrt{\frac{2k}{M}}$.

$\hbar \omega \frac{3}{2}$ ↑ ORBITAL W.F. ANTISYMM. $\Rightarrow$ SPIN SYMM. ($s=1$)
 $\hbar \omega \frac{1}{2}$ ↑ ORBITAL W.F. SYMMETRIC $\Rightarrow$ SPIN ANTISYMM ($s=0$)
 0

thus: the ground state has energy $\frac{1}{2} \hbar \omega$
 and $s=0 \Rightarrow$ no degeneracy.
 (deg = 1).

The 1st excited state $\frac{3}{2} \hbar \omega$
 has $s=1 \Rightarrow$ 3 possible spin m_s .
 (deg = 3)

Q2: Write the states as
 $|n\rangle \otimes \chi_{s,m}$ where $|n\rangle$ refers
 to the harmonic oscillator and
 $\chi_{s,m}$ ($s=0$ or 1) is the spin
 of the two particle system.

$$H' = \alpha (x_1 - x_2)^4 \hat{S}_1 \cdot \hat{S}_2 = \alpha x^4 \frac{\hbar^2}{2} \left(s(s+1) - \frac{3}{4} - \frac{3}{4} \right)$$

$$= \begin{cases} -\frac{3\alpha\hbar^2}{4} x^4 & \text{for } s=0 \text{ (ground state)} \\ +\frac{\alpha\hbar^2}{4} x^4 & \text{for } s=1 \text{ (1st excited st.)} \end{cases}$$

$$\langle 0 | x^4 | 0 \rangle = \left(\frac{\hbar}{2m\omega} \right)^2 \langle 0 | (a + a^\dagger)^4 | 0 \rangle = \frac{3\hbar^2}{4m^2\omega^2}$$

$$\langle 0 | (a + a^\dagger)(a + a^\dagger)(a + a^\dagger)(a + a^\dagger) | 0 \rangle =$$

$$= \langle 0 | a (a^2 + a^\dagger a + a a^\dagger + a^{\dagger 2}) a^\dagger | 0 \rangle =$$

$$= \langle 0 | a a^\dagger a a^\dagger + a^2 a^{\dagger 2} | 0 \rangle = 1 + 2 = 3 //$$

$$\langle 1 | x^4 | 1 \rangle = \left(\frac{\hbar}{2m\omega} \right)^2 \langle 1 | (a + a^\dagger)^4 | 1 \rangle = \frac{15\hbar^2}{4m^2\omega^2}$$

$$\langle 1 | (a + a^\dagger)(a + a^\dagger)(a + a^\dagger)(a + a^\dagger) | 1 \rangle =$$

$$= \langle 0 | a (a^2 a^{\dagger 2} + a a^\dagger a a^\dagger + a^\dagger a^2 a^\dagger + a a^\dagger a^\dagger a + a^\dagger a a^\dagger a + a^{\dagger 2} a^2) a^\dagger | 0 \rangle =$$

$$6 + 2 + 2 + 2 + 1 = 15 //$$

Hence:

$$E_0^{(1)} = - \frac{3\alpha\hbar^2}{4} \times \frac{3\hbar^2}{4m^2\omega^2} = - \frac{9\alpha\hbar^4}{8M\kappa}$$

$$E_1^{(1)} = + \frac{\alpha\hbar^2}{4} \cdot \frac{15\hbar^2}{4m^2\omega^2} = \frac{15\alpha\hbar^4}{8M\kappa}$$

PROBLEM 3

Let $\mathcal{H}_{s=1}^{(1)}$, $\mathcal{H}_{s=1}^{(2)}$, $\mathcal{H}_{s=1}^{(3)}$ be the part. 1, 2 and 3 Hilbert spaces.

Q1 Take the first 2 part's and combine them:

$$\mathcal{H}_{s=1}^{(1)} \otimes \mathcal{H}_{s=1}^{(2)} = \mathcal{H}_{s=0}^{(1,2)} \oplus \mathcal{H}_{s=1}^{(1,2)} \oplus \mathcal{H}_{s=2}^{(1,2)}$$

meaning that two $s=1$ part's combine to give $s=0$, 1 or 2.

Now take the third part and combine it with the previous ones

$$\mathcal{H}_{s=1}^{(3)} \otimes \mathcal{H}_{s=0}^{(1,2)} = \mathcal{H}_{s=1}^{(1,2,3)}$$

$$\mathcal{H}_{s=1}^{(3)} \otimes \mathcal{H}_{s=1}^{(1,2)} = \mathcal{H}_{s=0}^{(1,2,3)} \oplus \mathcal{H}_{s=1}^{(1,2,3)} \oplus \mathcal{H}_{s=2}^{(1,2,3)}$$

$$\mathcal{H}_{s=1}^{(3)} \otimes \mathcal{H}_{s=2}^{(1,2)} = \mathcal{H}_{s=1}^{(1,2,3)} \oplus \mathcal{H}_{s=2}^{(1,2,3)} \oplus \mathcal{H}_{s=3}^{(1,2,3)}$$

where $\mathcal{H}_{S_{TOT}}^{(1,2,3)}$ is the total space.

All together:

	deg. (= how many indep. states)
$S_{TOT} = 0$	1 (means No DEG).
1	3
2	2
3	1

check: $1 \times 1 + 3 \times 3 + 2 \times 2 + 1 \times 1 = 27 = 3^3$.

Q2: They are Bosons $\Rightarrow$ TOTAL STATE must be symmetric $\Rightarrow$ spin ANTISYMM.
 $\Rightarrow$ no value of $m = 0, \pm 1$ can be repeated.

$$\Rightarrow \chi^{spin} = \frac{1}{\sqrt{6}} \left(|1+\rangle |0\rangle \chi^- - |1+\rangle |1-\rangle |0\rangle - |0\rangle |1+\rangle |1-\rangle + |0\rangle |1-\rangle |1+\rangle + |1-\rangle |1+\rangle |0\rangle - |1-\rangle |0\rangle |1+\rangle \right)$$

the $\frac{1}{\sqrt{6}}$ is so that $\|\chi\|^2 = 1$ if the $|m\rangle$ states are properly normalized,

PROBLEM 4

$$\psi_-(x) = e^{ipx} + a e^{-ipx}$$

$$\psi_+(x) = b e^{ipx}$$

$$E = \frac{p^2}{2m}$$

set $\hbar = 1$
for convenience

continuity at $x=0 \Rightarrow 1+a=b$

$$\int_{-\epsilon}^{+\epsilon} \left(-\frac{1}{2m} \psi''(x) + \lambda \delta(x) \psi(x) \right) dx = E \int_{-\epsilon}^{+\epsilon} \psi(x) dx$$

$$\Downarrow$$

$$-\frac{1}{2m} \psi'(x) \Big|_{-\epsilon}^{+\epsilon} + \lambda \psi(0) = E \cdot 0$$

$$\Rightarrow -\frac{1}{2m} (ipb - ip + ipa) + \lambda(1+a) = 0$$

$$\Rightarrow a = \frac{-\lambda}{\lambda - ip/m}, \quad b = \frac{-ip/m}{\lambda - ip/m}$$

$$\Rightarrow R = |a|^2 = \frac{\lambda^2}{\lambda^2 + \frac{p^2}{m^2}}$$

$$T = |b|^2 = \frac{\frac{p^2}{m^2}}{\lambda^2 + \frac{p^2}{m^2}} = 1 - R \quad \checkmark$$

PROBLEM B

$$\text{Using } L_x = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$L_y = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & i \\ 0 & i & 0 \end{pmatrix}$$

$$L_z = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

(Can be obtained with $L_{\pm}$ if you forget).

$$H = \frac{\hbar\omega}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & -1 \\ 0 & -1 & 0 \end{pmatrix}$$

$$\text{Q1: } \begin{vmatrix} -\lambda & 1 & 0 \\ 1 & -\lambda & -1 \\ 0 & -1 & -\lambda \end{vmatrix} = 0 \quad \lambda_1 = 0, \lambda_2 = \sqrt{2}, \lambda_3 = -\sqrt{2}$$

$\Rightarrow$ energies: $0, \pm \hbar\omega$.

Q2: Energy eigenstates:

$$|E=0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} \quad |E=\hbar\omega\rangle = \frac{1}{2} \begin{pmatrix} 1 \\ \sqrt{2} \\ -1 \end{pmatrix}$$

$$|E=-\hbar\omega\rangle = \frac{1}{2} \begin{pmatrix} -1 \\ \sqrt{2} \\ 1 \end{pmatrix}$$

At $t = 0$:

$$|\psi, t=0\rangle = \frac{1}{\sqrt{2}}(|+\rangle - |-\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} \equiv |\psi_0\rangle =$$

$$= |E=0\rangle \langle E=0 | \psi_0 \rangle + |E=\hbar\omega\rangle \langle E=\hbar\omega | \psi_0 \rangle + \\ + |E=-\hbar\omega\rangle \langle E=-\hbar\omega | \psi_0 \rangle =$$

$$= \frac{1}{\sqrt{2}} |E=+\hbar\omega\rangle - \frac{1}{\sqrt{2}} |E=-\hbar\omega\rangle$$

At time $t > 0$:

$$|\psi, t\rangle = \frac{1}{\sqrt{2}} \left(e^{-i\omega t} |E=+\hbar\omega\rangle - e^{+i\omega t} |E=-\hbar\omega\rangle \right)$$

$$= \frac{1}{2\sqrt{2}} \begin{pmatrix} e^{-i\omega t} & e^{i\omega t} \\ \sqrt{2} (e^{-i\omega t} - e^{i\omega t}) \\ -e^{-i\omega t} - e^{i\omega t} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \cos \omega t \\ -\sqrt{2} i \sin \omega t \\ -\cos \omega t \end{pmatrix}$$

PROBLEM 6 set $\hbar=1$ for convenience.

Let $\psi_1(x) = e^{-a|x|}$ $\psi_2(x) = e^{-ax^2/2}$.

Q1: Be careful that ψ_1'' has a $\delta(x)$ in it.

Better to write:

$$\int_{-\infty}^{+\infty} \psi_1 \psi_1'' dx = - \int_{-\infty}^{+\infty} \psi_1'^2 dx = -2 \int_0^{\infty} \psi_1'^2 dx = -a$$

Also $\int_{-\infty}^{\infty} |x| \psi_1^2 dx = 2 \int_0^{\infty} x \psi_1^2 dx = \frac{1}{2a^2}$.

and $\int_{-\infty}^{+\infty} \psi_1^2 dx = 2 \int_0^{\infty} \psi_1^2 dx = \frac{1}{a}$

$$\bar{H}_1 = \frac{-\frac{1}{2m}(-a) + k \cdot \frac{1}{2a^2}}{1/a} = \frac{a^2}{2m} + \frac{k}{2a}$$

minimum for $a = \left(\frac{km}{2}\right)^{1/3}$

$$\Rightarrow \bar{H}_{1, \min} = \left(\frac{2}{mk}\right)^{1/3} \cdot \frac{3}{4} k \approx 0.945 \cdot \left(\frac{k^2}{m}\right)^{1/3}$$

Similarly:

$$\langle \psi_2 | \psi_2 \rangle = \int_{-\infty}^{+\infty} \psi_2^2 dx = \sqrt{\frac{\pi}{a}}$$

$$\int_{-\infty}^{+\infty} \psi_2 \psi_2'' dx = - \int_{-\infty}^{+\infty} \psi_2'^2 dx = - \frac{\sqrt{\pi a}}{2}$$

$$\int_{-\infty}^{+\infty} |x| \psi_2^2 dx = 2 \cdot \int_0^{\infty} x \psi_2^2 dx = \frac{1}{a}$$

$$\Rightarrow \bar{H}_2 = \frac{\frac{\sqrt{\pi a}}{4m} + \frac{k}{a}}{\sqrt{\pi/a}} = \frac{a}{4m} + \frac{k}{\sqrt{\pi a}}$$

$$\text{min for } a = \left(\frac{2km}{\sqrt{\pi}} \right)^{2/3}$$

$$\Rightarrow \bar{H}_{2\text{min}} = \frac{3}{2} \left(\frac{k^2}{2\pi m} \right)^{1/3} \approx 0.813 \cdot \left(\frac{k^2}{m} \right)^{1/3}$$

Q2: The second value is better because it is lower, and the variational approach always gives an upper bound.